Closure Temperature

R. Abart, E. Petrishcheva

Free University Berlin

21 April 2006



- 1 Physical background;
- 2 Transformation of time variable;
- 3 Atomistic interpretation of "closing", "freezing";



Layman asks: What is closing temperature???

Expert answers: Closing temperature is an individual concept !!!



The closure temperature of a geochronolgical system may be defined as the temperature at the time corresponding to the apparent age. ...

The temperature recorded by a "frozen" chemical system, in which a solid phase in contact with a large reservoir has cooled slowly from high temperatures, is formally identical with geochronological closing temperature

M.H. Dodson (1973),

Contrib.Min.Pet., 40, 259-274



Diffusion coefficient usually follows the Arrhenius relation

$$D(T)=D_{\infty}e^{\frac{-E}{RT}},$$

where D_{∞} is the diffusion coefficient for RT >> E, i.e. for infinitely high temperature. If temperature varies with time, T = T(t):

$$D(t) = D_{\infty} e^{\frac{-E}{RT(t)}}$$

and the diffusion equation becomes

$$\frac{\partial c}{\partial t} = D(t) \frac{\partial^2 c}{\partial x^2}$$



Transformation of the time variable

Use $T_0 = T(t = 0)$, then:

$$D(t) = D_0 \exp\left[-rac{E}{R}\left(rac{1}{T(t)} - rac{1}{T_0}
ight)
ight],$$

where D_0 is D at T_0 . Introduce $\zeta(t)$ to scale D to the correct value at time t

$$D(t) = D_0 \zeta(t)$$

Transformed (compressed) time is then defined as:

$$t'=\int_0^t\zeta(t)\ dt$$

or replacing for $\zeta(t)$

$$t' = \int_0^t \exp\left[-\frac{E}{R}\left(\frac{1}{T(t)} - \frac{1}{T_0}\right)\right] dt$$
Freie Universität

by differentiation we obtain

 $dt' = \zeta(t) dt$

The diffusion equation then becomes

$$\frac{\partial c}{\partial t} = D_0 \zeta(t) \frac{\partial^2 c}{\partial x^2},$$

dividing both sides by $\zeta(t)$ and using $dt' = \zeta(t) dt$ yields:

$$\frac{\partial c}{\partial t'} = D_0 \frac{\partial^2 c}{\partial x^2},$$

i.e. the problem is reduced to a *standard* diffusion problem with constant diffusion coefficient.

Berlin

Cooling History

Let T decrease monotonously from T_0 at a constat cooling rate s so that:

$$T(t) = \frac{T_0}{1 + (st/T_0)}$$

The time dependence of D may then be expressed as:

$$D(t) = D_0 \exp\left[-rac{Est}{RT_0^2}
ight]$$

Let $\gamma = Es/RT_0^2$ then this becomes

$$D(t)=D_0e^{-\gamma t},$$

and

$$t'=rac{1}{\gamma}\left(1-e^{-\gamma t}
ight).$$

t' can not increase to infinity but is limited to a finite value,

$$\lim_{t\to\infty}(t')=1/\gamma.$$



For a small time interval $d\tau$ the mean squared distance traveled by a particle in a small time interval $d\tau$ is $2D_{\tau}d\tau$ (Einstein equation). The total mean squared distance traveled after time t is

$$\left(X^2\right) = \int_0^t 2D_\tau d\tau$$

given cooling at a constant cooling rate and inserting for $D_{ au}$ yields:

$$(X^2) = 2D_0 \frac{1}{\gamma} \left(1 - e^{-\gamma t}\right),$$

The limiting distance traveled by a particle after $t \to \infty$ is then:

$$\left(X^{2}\right) = 2D_{0}\frac{1}{\gamma} = 2D_{\infty}e^{-\frac{E}{RT_{0}}}\frac{RT_{0}^{2}}{Es}$$



Let *a* be a typical diffusion distance in a geological process , then the closuretemperature, T_c , may be obtained by finding the minimum initial temperature, T_0 , which gives a diffusion distance of *a* :

$$a^{2} = \frac{2D_{\infty}RT_{c}^{2}}{Es}e^{-\frac{E}{RT_{c}}}$$

If $T_0 > T_c$, then the mean squared distance traveled by a particle $> a^2$, if $T_0 < T_c$, then the mean squared distance traveled $< a^2$ and the mineral grain acts as a *closed system*. Rearrangement gives:

$$T_c = \frac{E/R}{\ln\left[2D_{\infty}RT_c^2/a^2Es\right]}$$

This may be iterated to obtain T_c .

